Lecture 10 part 2

Slice method

1. Particle method (Lect. 8-9):

assumption of particle motion within motionless environment.

2. Basic approaches of <u>slice method</u>: 1938-1939 (Bjerknes J., Petersen S., Shishkin N.)

Let's pick out a sufficiently large layer of air

S' cross section(inner)



Slice method assumptions:

- all variations of the quantities within the picked out layer come out adiabatically;
- no advection (horizontal mass motion) is observed;
- the mass of air above any layer remains unchanged.

The latter means that fluxes of air mass through section S['] and S'' are equal:

$$\rho'S'w' = -\rho''S''w''$$
 (10.8)

 ρ is density

$$\rho' \approx \rho''$$

$$S'w' = -S''w''$$
 (10.9)

To derive the stability criterion:

 ΔT is temperature difference, T' is updraft flux temp., T'' is the downdraft one at z level $\Delta T = T' - T''$

 $\Delta T > 0$ unstable state(positive acceleration)

 $\Delta T < 0$ stable state(negative acceleration)

Suppose, particles come from lower level to upper one:

$$T' = T_1 - \gamma'(z - z_1)$$
 (10.10)
 $T'' = T_2 + \gamma''(z_2 - z_1)$ (10.11)

 γ' is adiabatic lapse rate in updraft flux γ'' is adiabatic lapse rate in downdraft flux

$$z - z_1 = w' \Delta t$$
 (10.12)
 $z_2 - z = -w'' \Delta t$ (10.13)

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From (10.10)-(10.13) we obtain:

$$\Delta T = T_1 - T_2 - (\gamma' w' - \gamma'' w'') \Delta t \qquad (10.14)$$

$$T_1 - T_2 = \gamma \ (z_2 - z_1) = \gamma \ (z_2 - z) + \gamma \ (z - z_1) = \gamma \ \Delta t (w' - w'')$$

$$\Delta T = \left[\gamma'(w' - w'') - (\gamma'w' - \gamma''w'') \right] \Delta t$$
 (10.15)

Let's use (10.9)

$$S'w' = -S''w'' (10.9)$$

$$\Delta T = \left[(\gamma'(w' - w'') - (\gamma'w' - \gamma''w'') \right] \Delta t$$
(10.15)

$$\Delta T = \left[(\gamma - \gamma' \quad) + \frac{S'}{S''} (\gamma - \gamma'') \right] w' \Delta t \qquad (10.16)$$

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Let's consider 3 cases:

1.
$$\gamma' = \gamma'' = \gamma_a$$

Updraft and downdraft air isn't saturated

$$\Delta T = (\gamma - \gamma_a)(1 + \frac{S'}{S''})w'\Delta t \qquad (10.17)$$

 Criteria of stability of both methods (parcel and slice ones) give the same result.



2.
$$\gamma' = \gamma'' = \gamma'_a$$

 Updraft and downdraft air fluxes are saturated

$$\Delta T = (\gamma - \gamma_a')(1 + \frac{S'}{S''})w'\Delta t$$

• Criteria of stability of both methods (parcel and slice ones) give the same result.



3.
$$\gamma' = \gamma'_a, \gamma'' = \gamma_a$$

• Updraft air is saturated and downdraft flux is not.

$$\Delta T = \left[(\gamma - \gamma_a') + \frac{S'}{S''} (\gamma - \gamma_a) \right] w' \Delta t \qquad (10.18)$$

• Cu clouds development

- From (10.18) $\Delta T > 0$ at $\gamma > \gamma_a$ and
- $\Delta T < 0 \quad at \ \gamma < \gamma \ a$ • Stability criteria of both methods are coincide if the stratification is absolutely unstable and absolutely stable. However, often

$$\gamma_a < \gamma < \gamma_a$$

$$\Delta T \leftarrow on \gamma but also S' / S''$$

Let's introduce a critical lapse rate

$$\gamma_{cr} \quad \alpha t \quad \Delta T = 0$$

$$\left[(\gamma - \gamma_a') + \frac{S'}{S''} (\gamma - \gamma_a) \right] w' \Delta t = 0$$

$$w' \neq 0, \quad \Delta t \neq 0$$

$$\left[(\gamma - \gamma_a') + \frac{S'}{S''} (\gamma - \gamma_a) \right] = 0$$

$$\gamma_{cr} = \frac{\gamma'_a + \gamma_a (S' / S'')}{1 + S' / S''} \quad (10.19)$$

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$$\gamma_{cr} \rightarrow \gamma_{a}^{'} \text{ at small values of } \left(\frac{S'}{S''}\right)$$

$$\gamma_{cr} \rightarrow \gamma_{a} \text{ at } l \arg e \text{ values of } \left(\frac{S'}{S''}\right)$$

$$\gamma > \gamma_{cr} \ (\Delta T > 0) \qquad \gamma < \gamma_{cr} \ (\Delta T < 0)$$

$$\frac{S'}{S''} < \frac{\gamma - \gamma'_{a}}{\gamma_{a} - \gamma} \text{ at } \Delta T > 0 \qquad (10.20)$$

$$\frac{S'}{S''} < \frac{\gamma - \gamma'_a}{\gamma_a - \gamma} \quad at \quad \Delta T < 0 \tag{10.21}$$

 at conditionally unstable stratification the atmosphere is unstable for small size particles and it's stable for LARGE size particles =>>the atmospere is said to be <u>selectively unstable</u>