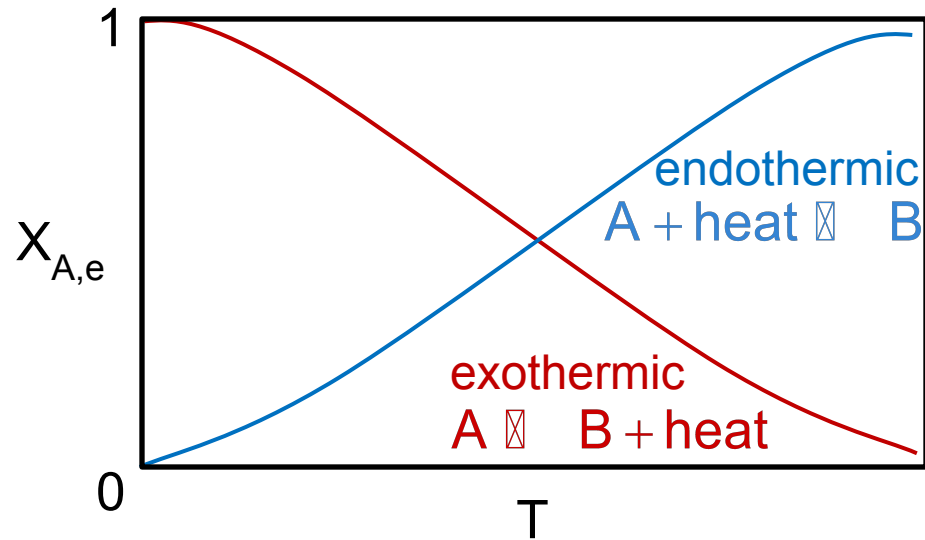


Review: Equilibrium Conversion X_{Ae}



Example) $A \rightleftharpoons B$ $C_{A0}=1$ $C_{B0}=0$

$$K_C = \frac{C_{Be}}{C_{Ae}} = \frac{\cancel{C_{A0}}(0 + X_{Ae})}{\cancel{C_{A0}}(1 - X_{Ae})} \rightarrow K_C = \frac{X_{Ae}}{1 - X_{Ae}} \quad \text{Rearrange to solve in terms of } X_{Ae}$$

$$\rightarrow K_C(1 - X_{Ae}) = X_{Ae} \rightarrow K_C = X_{Ae} + K_C X_{Ae} \rightarrow K_C = X_{Ae}(1 + K_C)$$

$$\rightarrow \boxed{\frac{K_C}{(1 + K_C)} = X_{Ae}} \quad \text{This equation enables us to express } X_{ae} \text{ as a function of } T$$

Review: X_{Ae} and Temperature

$$X_{Ae} = \frac{1}{\frac{1}{K_C(T_2)} \exp\left[\frac{\Delta H_{RX}^\ddagger(T_R)}{R}\left(\frac{1}{T} - \frac{1}{T_2}\right)\right] + 1}$$

Clicker question material

Exothermic & $\Delta C_p = 0$: $\Delta H_{RX}^\ddagger(T_R) < 0$, when $T \uparrow$ $\exp\left[\frac{\Delta H_{RX}^\ddagger(T_R)}{R}\left(\frac{1}{T} - \frac{1}{T_2}\right)\right] \uparrow$ & $X_{Ae} \downarrow$

Makes sense from Le Chatelier's principle $A \rightleftharpoons B + \text{heat}$

Exothermic rxn produces heat \rightarrow

increasing temp adds heat (product) & pushes rxn to left (lower conversion)

Endothermic & $\Delta C_p \approx 0$: $\Delta H_{RX}^\ddagger(T_R) > 0$, when $T \uparrow$ $\exp\left[\frac{\Delta H_{RX}^\ddagger(T_R)}{R}\left(\frac{1}{T} - \frac{1}{T_2}\right)\right] \downarrow$ & $X_{Ae} \uparrow$

Makes sense from Le Chatelier's principle $A + \text{heat} \rightleftharpoons B$

Heat is a reactant in an endothermic rxn \rightarrow

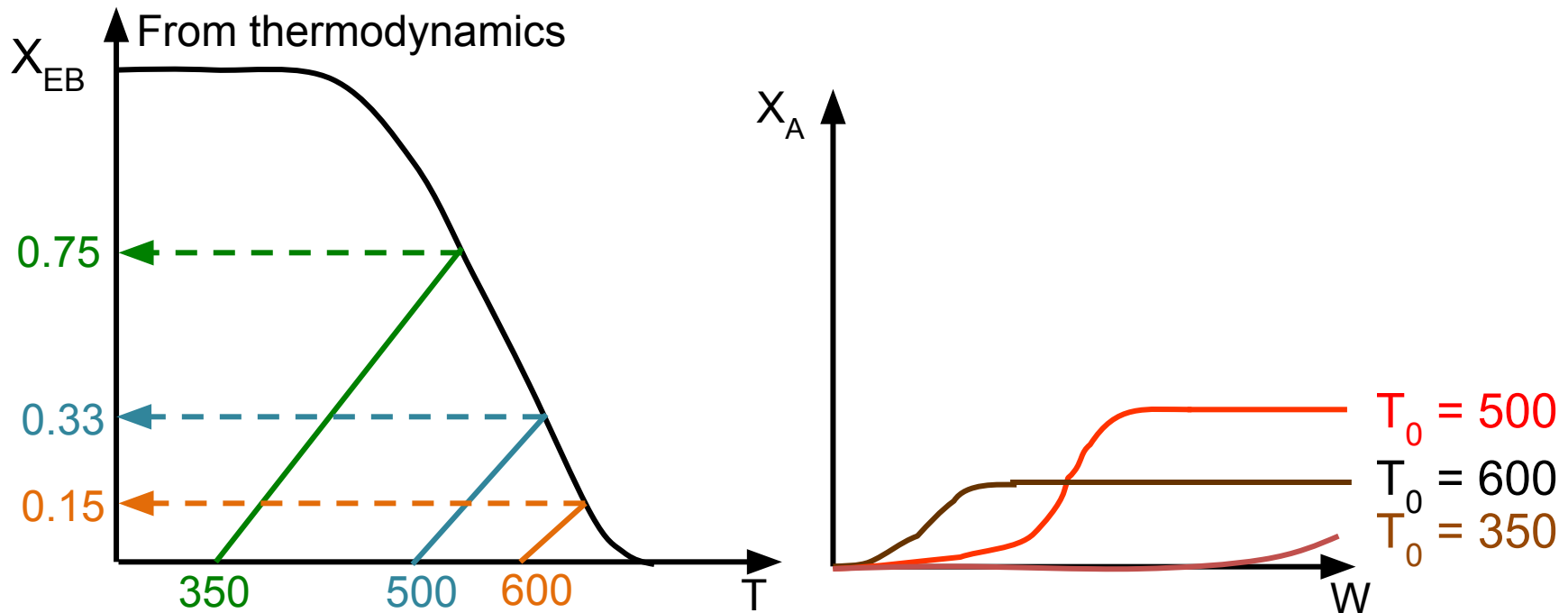
increasing temp adds reactant (heat) & pushes rxn to right (higher conversion)

Review: Optimum Feed Temperature

For reversible, exothermic rxns, optimize feed temperature to maximize X_A

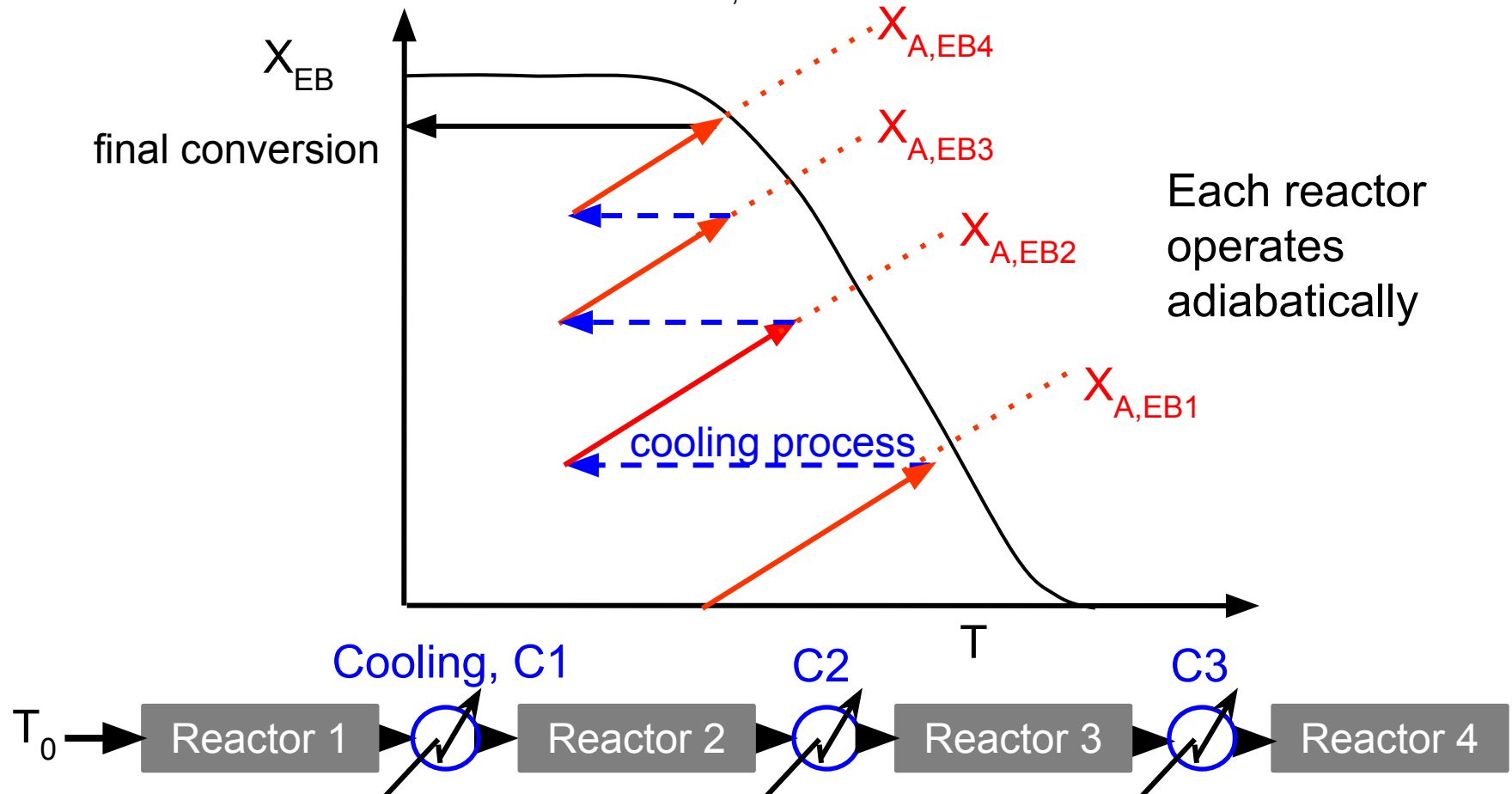
High T_0 : moves $X_{A,EB}$ line to the right. Rxn reaches equilibrium fast, but low X_A

Low T_0 would give high $X_{A,e}$ but the specific reaction rate k is so small that most of the reactant passes through the reactor without reacting (never reach $X_{A,e}$)



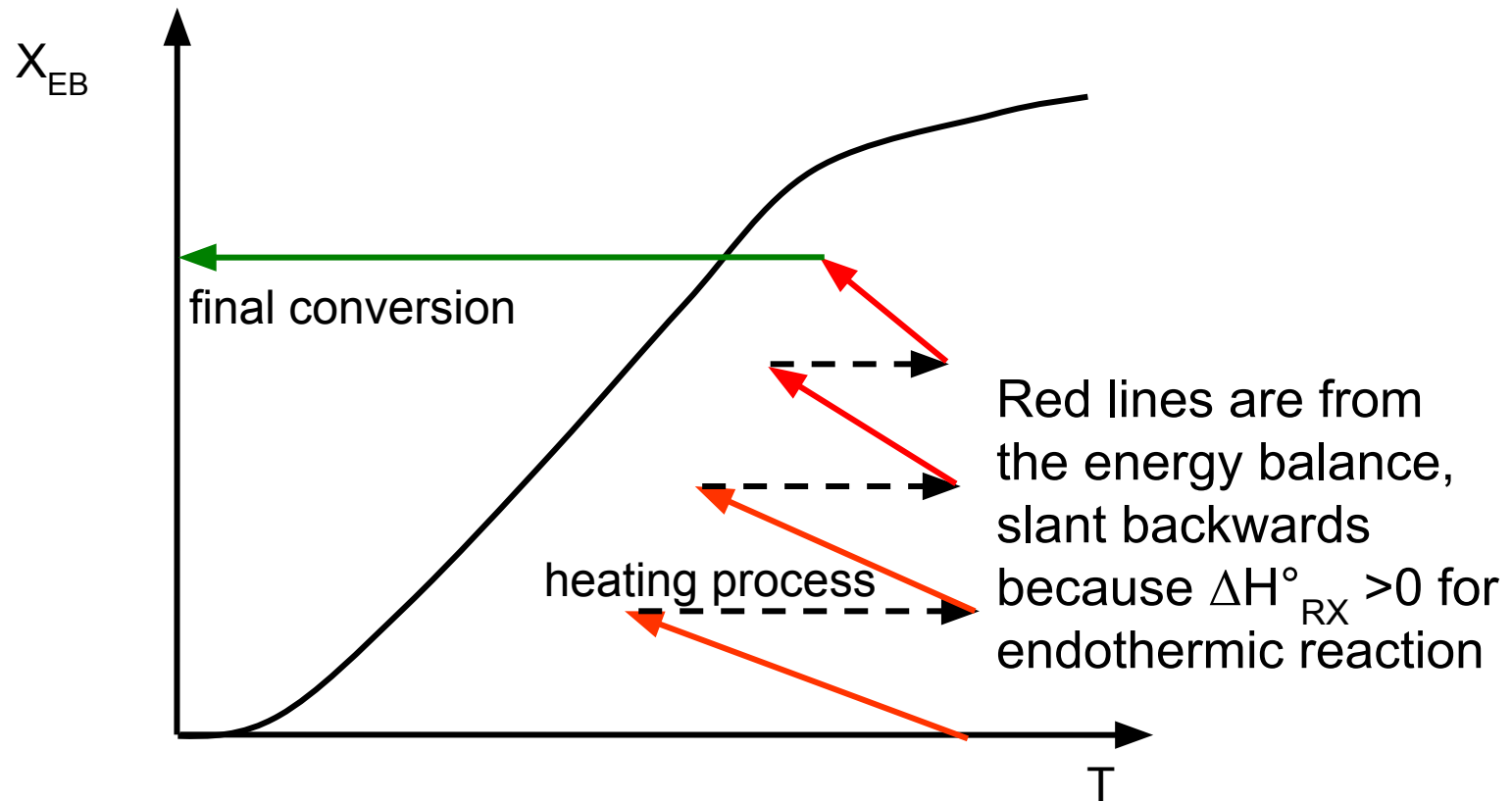
Review: Interstage Cooling

- Adiabatic operation of each reactor simplifies the energy balance
- Higher feed temp- reaction reaches equilibrium quickly but $X_{A,e}$ is low
- Lower feed temp- higher $X_{A,e}$ but reaction rate is too slow to be practical
- Cooling between reactors shifts $X_{A,EB}$ line to the left, increasing X_A



Review: Endothermic Reactions

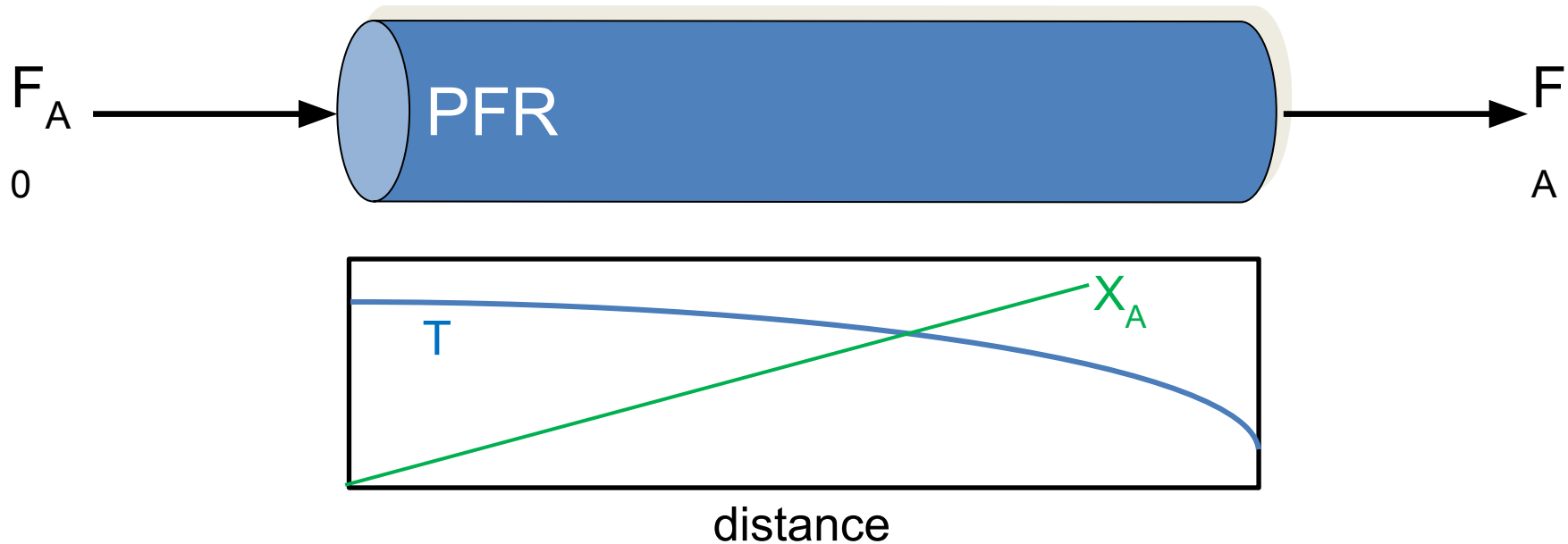
The equilibrium conversion increases with increasing temperature, so use interstage heating to increase the conversion



L14: Nonadiabatic PFR/PBR Operation and Reactor Stability

1. T changes with distance down reactor- differential form of EB must be used
2. Multiple steady states: more than one set of conditions satisfies both the energy balance & mole balance

Review: Application to a SS PFR



Negligible shaft work ($\dot{W}_s=0$) and **adiabatic** ($\dot{Q}=0$)

- Use TEB to construct a table of T as a function of X_A
- Use $k = Ae^{-E/RT}$ to obtain k as a function of X_A
- Use stoichiometry to obtain $-r_A$ as a function of X_A
- Calculate:

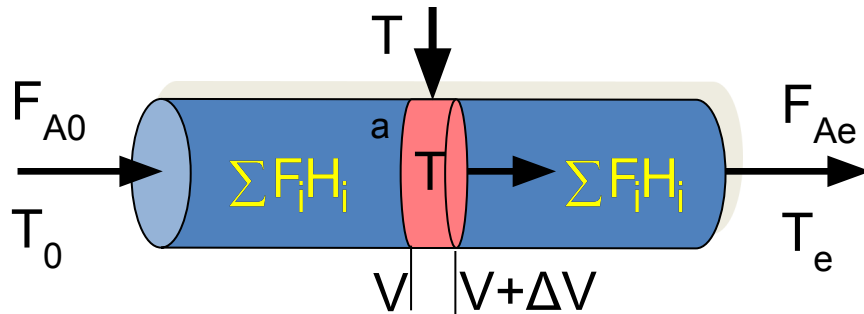
$$V = F_{A0} \int_{X_{A0}}^{X_A} \frac{dX_A}{-r_A(X_A, T)} \quad \text{may use numerical methods}$$

Steady-State PFR/PBR w/ Heat Exchanger

Heat is added or removed through the cylindrical walls of the reactor

$$\dot{Q} = U\Delta A(T_a - T) = Ua(T_a - T)\Delta V$$

$$a = \frac{A}{V}$$



Heat exchange area per volume of reactor

Energy balance on small volume of SS PFR:

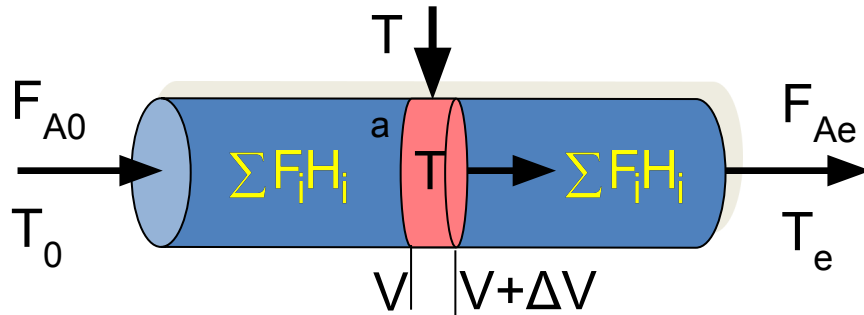
$$\dot{\Delta Q} - \dot{W}_s + \sum F_i H_i|_V - \sum F_i H_i|_{V+\Delta V} = 0$$

Plug in Q: $\rightarrow Ua(T_a - T)\Delta V - \dot{W}_s + \sum F_i H_i|_V - \sum F_i H_i|_{V+\Delta V} = 0$

Take limit as $\Delta V \rightarrow \infty$: $\rightarrow Ua(T_a - T) - \frac{d(\sum F_i H_i)}{dV} = 0$

Expand: $\rightarrow Ua(T_a - T) - \sum \frac{dF_i}{dV} H_i - \sum F_i \frac{dH_i}{dV} = 0$

TEB for PFR/PBR w/ Heat Exchanger



$$Ua(T_a - T) - \sum \frac{dF_i}{dV} H_i - \sum F_i \frac{dH_i}{dV} = 0$$

Substitute the differentials: $\frac{dF_i}{dV} = r_i = v_i(-r_A)$ and $\frac{dH_i}{dV} = C_{pi} \frac{dT}{dV}$

$$\rightarrow Ua(T_a - T) - \sum H_i v_i (-r_A) - \sum F_i C_{pi} \frac{dT}{dV} = 0 \quad \sum v_i H_i = \Delta H_{RX}$$

$$\rightarrow Ua(T_a - T) - \Delta H_{RX} (-r_A) - \sum F_i C_{pi} \frac{dT}{dV} = 0 \quad \text{Solve for } dT/dV:$$

$$\rightarrow -\sum F_i C_{pi} \frac{dT}{dV} = \Delta H_{RX} (-r_A) - Ua(T_a - T) \rightarrow \frac{dT}{dV} = \frac{\Delta H_{RX} (r_A) + Ua(T_a - T)}{\sum F_i C_{pi}}$$

Energy Balance for Tubular Reactors

$$\frac{dT}{dV} = \frac{\overbrace{\Delta H_{RX}(r_A)}^{\text{Heat generated}} + \underbrace{Ua(T_a - T)}_{\text{Heat removed}}}{\sum F_i C_{Pi}}$$

Multiply Ua and $(T_a - T)$ by -1 ($-1 \times -1 = 1$) \rightarrow

$$\frac{dT}{dV} = \frac{\overbrace{\Delta H_{RX}(r_A)}^{\text{Heat generated}} - \underbrace{Ua(T - T_a)}_{\text{Heat removed}}}{\sum F_i C_{Pi}}$$

$$\left. \begin{aligned} \Delta H_{RX} &= \Delta H_{RX}^\ominus(T_R) + \Delta C_P(T - T_R) \\ F_i &= F_{A0}(\Theta_i + \nu_i X_A) \end{aligned} \right\} \begin{array}{l} \text{Substitute and multiply} \\ \text{out the denominator} \end{array}$$

Switched sign & order in bracket

$$\frac{dT}{dV} = \frac{\left[\Delta H_{RX}^\ominus(T_R) + \Delta C_P(T - T_R) \right] (r_A) + Ua(T_a - T)}{F_{A0} \sum (\Theta_i C_{Pi} + \underbrace{C_{Pi} \nu_i X_A}_{\text{for } A \rightarrow B})}$$

$$\nu_i C_{Pi} = \frac{b}{a} C_{PB} - C_{PA} = \Delta C_P$$

$$\rightarrow \frac{dT}{dV} = \frac{\left[\Delta H_{RX}^\ominus(T_R) + \Delta C_P(T - T_R) \right] (r_A) + Ua(T_a - T)}{F_{A0} (\sum \Theta_i C_{Pi} + \Delta C_P X_A)}$$

Energy balance for SS PFR, $\dot{W}_s = 0$

$$\frac{dX_A}{dV} = \frac{-r_A}{F_{A0}}$$

PFR energy balance is coupled to the PFR design eq, and PFR design eq is coupled to Arrhenius eq for k or K_{equil}

(these are the 3 equations that must be simultaneously solved)

Liquid Phase Reaction in PFR

A \boxtimes B liquid phase rxn carried out in PFR; $\dot{W}_S = 0$ & pure A enters the PFR

✓ Mole balance

$$\frac{dX_A}{dV} = \frac{-r_A}{F_{A0}}$$

✓ Rate law

$$-r_A = k \left(C_A - \frac{C_B}{K_C} \right) \quad \text{with}$$

$$k = k_1 \exp \left[\frac{E}{R} \left(\frac{1}{T_1} - \frac{1}{T} \right) \right]$$

$$K_C(T) = K_C(T_2) \exp \left[\frac{\Delta H_{RX}^\boxtimes(T_R)}{R} \left(\frac{1}{T_2} - \frac{1}{T} \right) \right]$$

✓ Stoichiometry

$$C_A = C_{A0}(1 - X_A) \quad C_B = C_{A0}X_A$$

✓ Combine

$$\frac{dX_A}{dV} = \frac{k[1 - X_A - X_A/K_C]}{v_0}$$

✓ Energy balance

$$\frac{dT}{dV} = \frac{[\Delta H_{RX}^\boxtimes(T_R) + \Delta C_P(T - T_R)](r_A) + Ua(T_a - T)}{F_{A0}(\sum \Theta_i C_{Pi} + \Delta C_P X_A)}$$

Solve these equations simultaneously with an ODE solver (Polymath)

If this were a gas phase rxn w/ pressure drop, change stoichiometry accordingly & include an equation for $d\Delta P/dW$

Review: Nonisothermal CSTR

Isothermal CSTR: feed temp = temperature inside the CSTR

Case 1: Given F_{A0} , C_{A0} , A , E , C_{pi} , H°_i , and X_A , calculate T & V

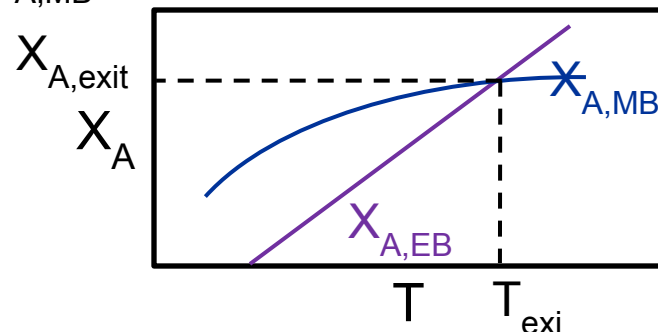
- Solve TEB for T at the exit ($T_{\text{exit}} = T_{\text{inside reactor}}$)
- Calculate $k = Ae^{-E/RT}$ where T was calculated in step a
- Plug the k calculated in step b into the design equation to calculate V_{CSTR}

Case 2: Given F_{A0} , C_{A0} , A , E , C_{pi} , H°_i , and V , calculate T & X_A

- Solve TEB for T as a function of X_A
- Solve CSTR design equation for X_A as a function of T (plug in $k = Ae^{-E/RT}$)
- Plot $X_{A,EB}$ vs T & $X_{A,MB}$ vs T on the same graph. The intersection of these 2 lines is the conditions (T and X_A) that satisfies the energy & mass balance

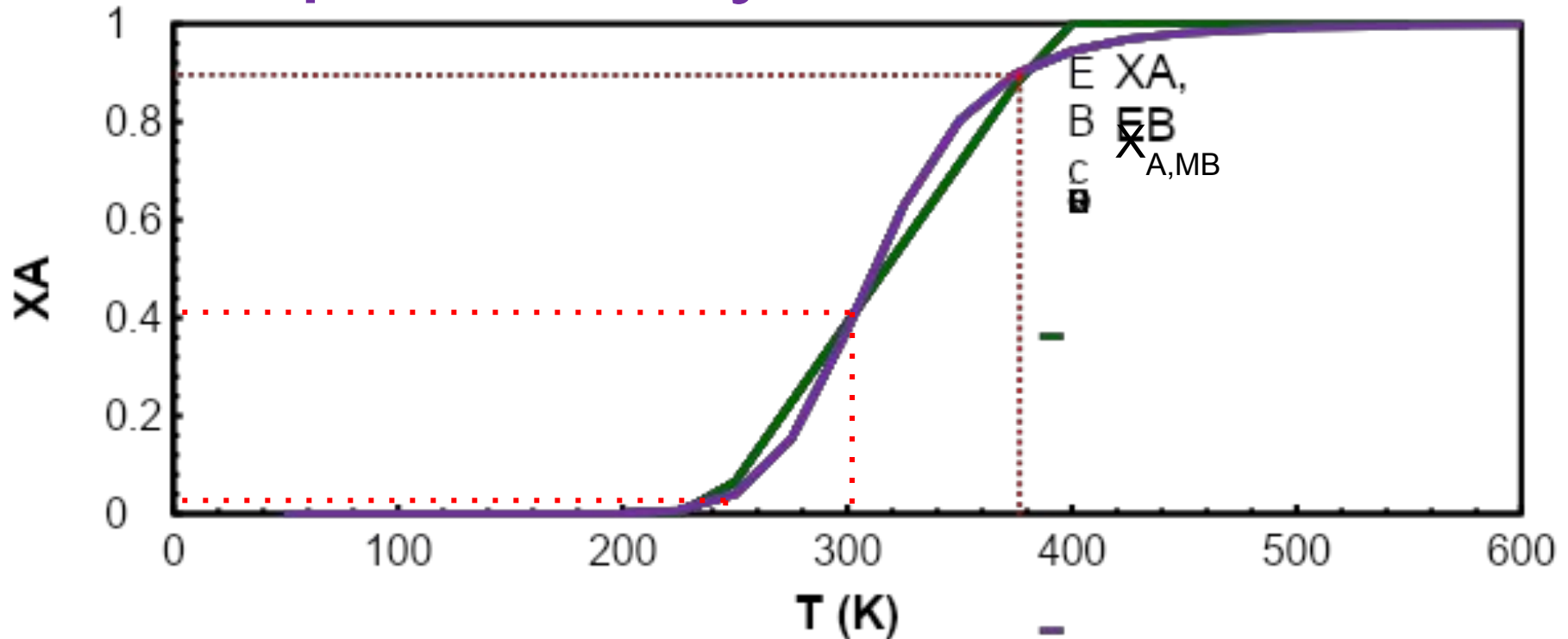
$X_{A,EB}$ = conversion determined from the TEB equation

$X_{A,MB}$ = conversion determined using the design equation



Intersection is T and X_A that satisfies both equations

Multiple Steady States in CSTR



- Plot of $X_{A,EB}$ vs T and $X_{A,MB}$ vs T
- Intersections are the T and X_A that satisfy both the mass balance and energy balance
- Multiple sets of conditions are possible for the same rxn in the same reactor with the same inlet conditions!

Reactor must operate near one of these steady states-
this requires knowledge of their stability!

Consider a jacketed CSTR with constant heat capacity, negligible shaft work, $\Delta C_p = 0$, first order kinetics, all feeds at the same temperature ($T_{i0} = T_0$), constant T_a in jacket, and an overall heat transfer coefficient

$$\dot{Q} = UA(T_a - T)$$

$$\text{TEB: } 0 = \dot{Q} - F_{A0} \sum_{i=1}^n \Theta_i C_{p,i} [T - T_{i0}] - \Delta H_{RX}(T) F_{A0} X_A$$

Substituting for $\dot{Q} = UA(T_a - T)$ and $\Delta H_{RX}(T) = \Delta H_{RX}^\ominus(T_R)$ since $\Delta C_p = 0$

$$\rightarrow 0 = UA(T_a - T) - F_{A0} \sum_{i=1}^n \Theta_i C_{p,i} [T - T_{i0}] - \Delta H_{RX}^\ominus(T_R) F_{A0} X_A$$

Bring terms that remove heat to other side of equation:

$$\rightarrow F_{A0} \sum_{i=1}^n \Theta_i C_{p,i} [T - T_{i0}] - UA(T_a - T) = -\Delta H_{RX}^\ominus(T_R) F_{A0} X_A$$

$$\rightarrow \underbrace{\sum_{i=1}^n \Theta_i C_{p,i} [T - T_{i0}] - \frac{UA(T_a - T)}{F_{A0}}}_{\text{Heat removed term} \equiv R(T)} = \underbrace{-\Delta H_{RX}^\ominus(T_R) X_A}_{\text{Heat generated term} \equiv G(T)}$$

A steady-state occurs when $R(T) = G(T)$

Even More Terms...

Heat removed term $\equiv R(T)$

Heat generated term $\equiv G(T)$

$$\overbrace{\sum_{i=1}^n \Theta_i C_{p,i} [T - T_{i0}] - \frac{UA(T_a - T)}{F_{A0}}} = \overbrace{-\Delta H_{RX}^{\boxtimes}(T_R) X_A}$$

$$C_{p0} = \sum \Theta_i C_{p,i} \quad V = \frac{F_{A0} X_A}{-r_A} \rightarrow \frac{-r_A V}{F_{A0}} = X_A$$

Substitute $\rightarrow C_{p0} [T - T_{i0}] - \frac{UA(T_a - T)}{F_{A0}} = -\Delta H_{RX}^{\boxtimes}(T_R) \left(\frac{-r_A V}{F_{A0}} \right)$

More substitutions: $\kappa = \frac{UA}{C_{p0} F_{A0}} \quad T_c = \frac{T_0 F_{A0} C_{p0} + UA T_a}{UA + C_{p0} F_{A0}} = \frac{\kappa T_a + T_0}{1 + \kappa}$

$$\Delta H_{RX}^{\boxtimes} = \Delta H_{RX}^{\boxtimes}(T_R)$$

$$\rightarrow C_{p0} (1 + \kappa) [T - T_c] = -\Delta H_{RX}^{\boxtimes} \left(\frac{-r_A V}{F_{A0}} \right)$$

Heat

removed: $R(T) = C_{p0} (1 + \kappa) [T - T_c]$

Heat

generated: $G(T) = -\Delta H_{RX}^{\boxtimes} \left(\frac{-r_A V}{F_{A0}} \right)$

Heat Removal Term and T_0

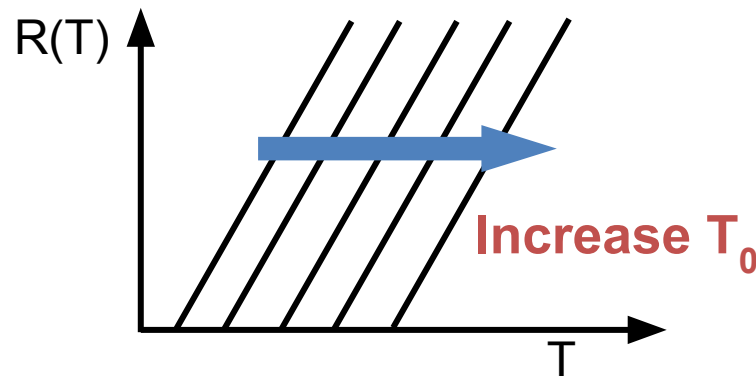
Heat removed: $R(T)$

Heat generated $G(T)$

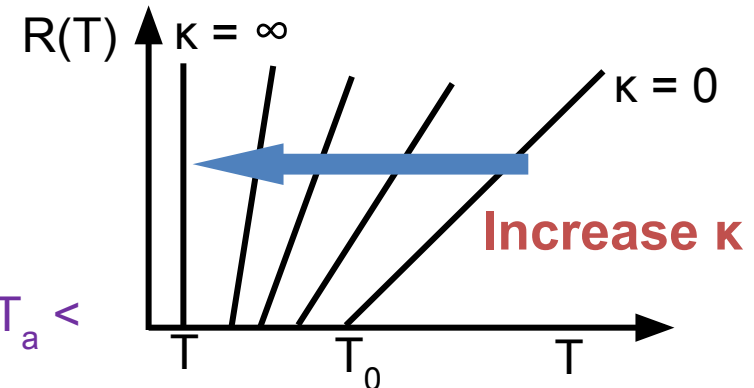
$$C_{p0}(1+\kappa)[T - T_C] = -\Delta H_{RX} \left(\frac{-r_A V}{F_{A0}} \right)$$

$$\kappa = UA/C_{p0}F_{A0} \quad T_C = \frac{\kappa T_a + T_0}{1 + \kappa}$$

$R(T)$ line has slope of $C_{p0}(1+\kappa)$



When T_0 increases, slope stays same & line shifts to right



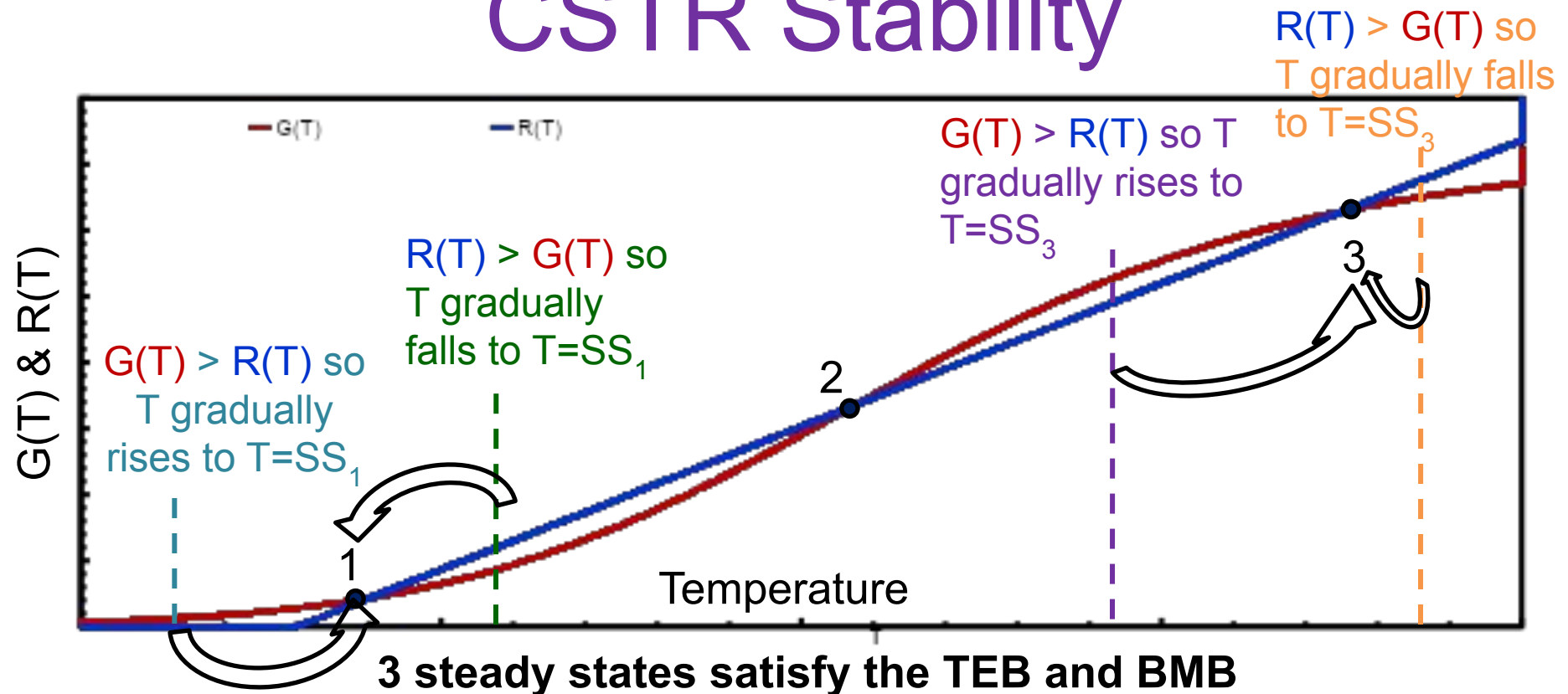
When κ increases from lowering F_{A0} or increasing heat exchange, slope and x-intercept moves

$T_a < T_0$: x-intercept shifts left as $\kappa \uparrow$

$T_a > T_0$: x-intercept shifts right as $\kappa \uparrow$

$\kappa=0$, then $T_C = T_0$ $\kappa=\infty$, then $T_C = T_a$

CSTR Stability

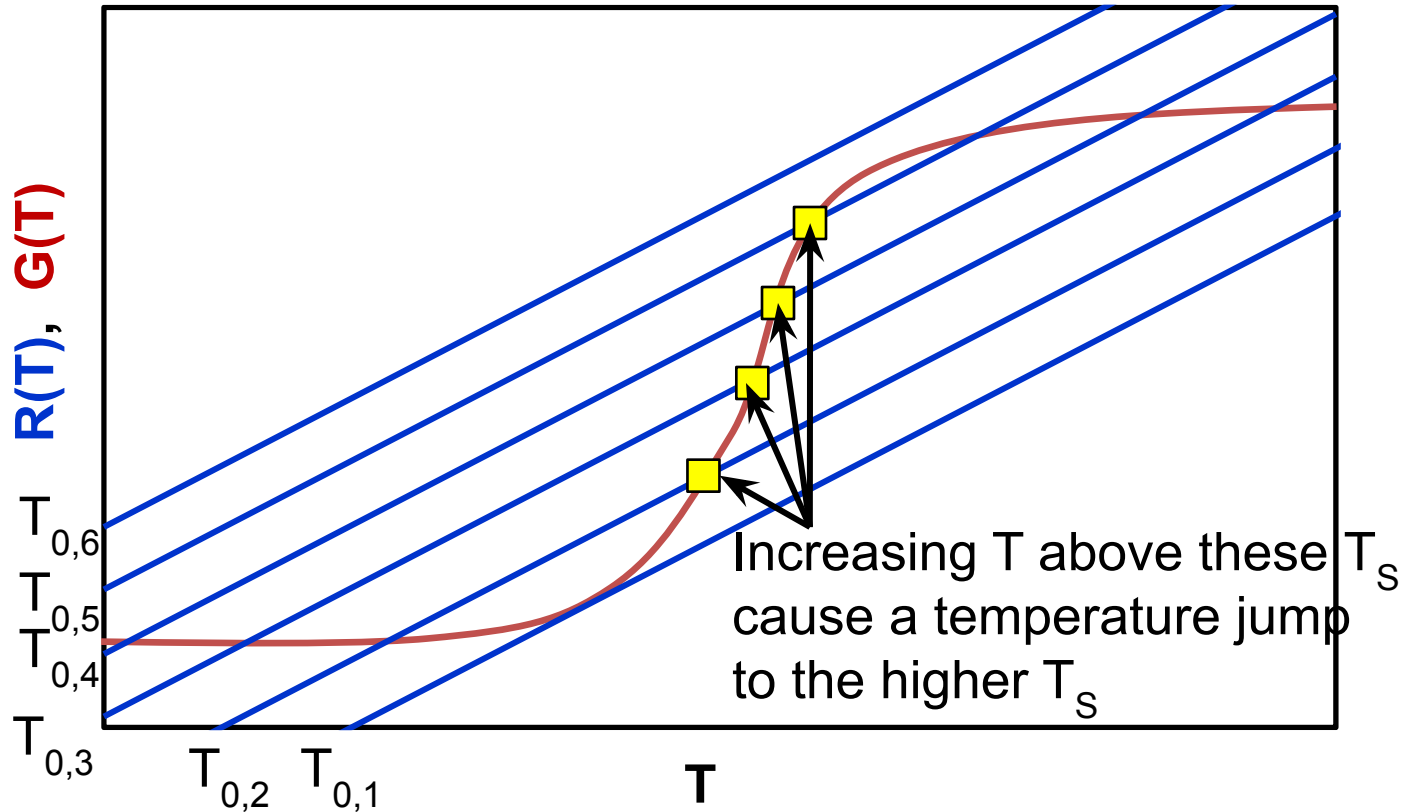


- Suppose a disturbance causes the reactor T to drift to a T between SS_1 & SS_2
- Suppose a disturbance causes the reactor T to drift to a T between SS_2 & SS_3
- Suppose a disturbance causes the reactor T to drop below SS_1
- Suppose a disturbance causes the reactor T to rise above SS_3

SS_1 and SS_3 are locally stable (return to them after temp pulse)

SS_2 is an unstable- do not return to SS_2 if there is a temp pulse

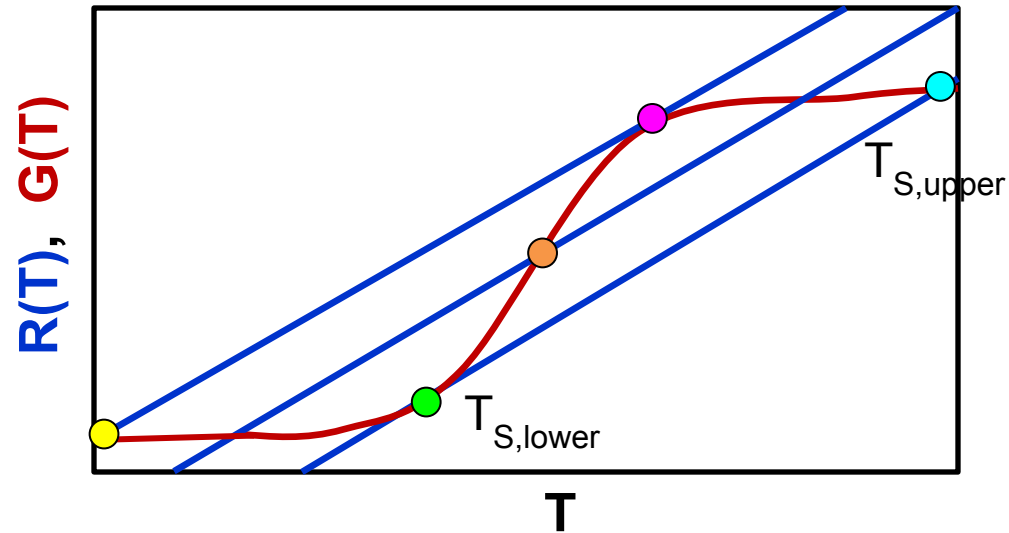
Multiple Steady States and T_0



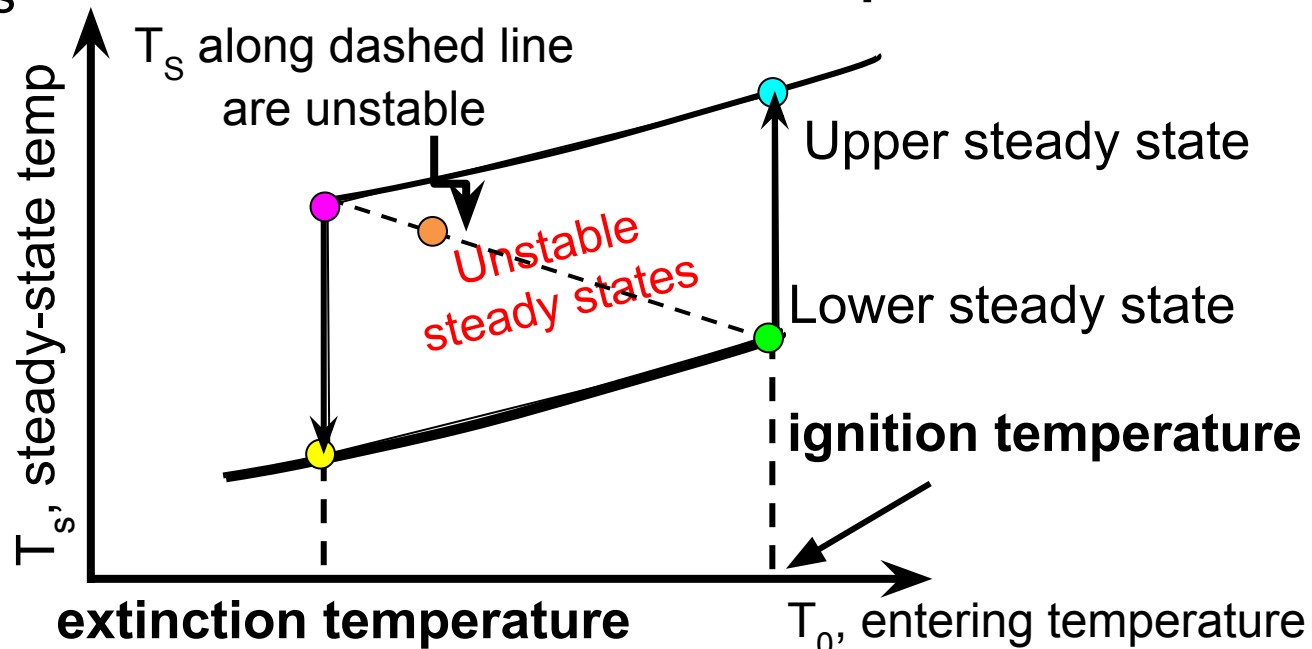
- Changing the inlet T will shift the steady state temperature (T_S)
- Notice that the number of steady state temperatures depends on T_0

Temperature Ignition-Extinction Curve

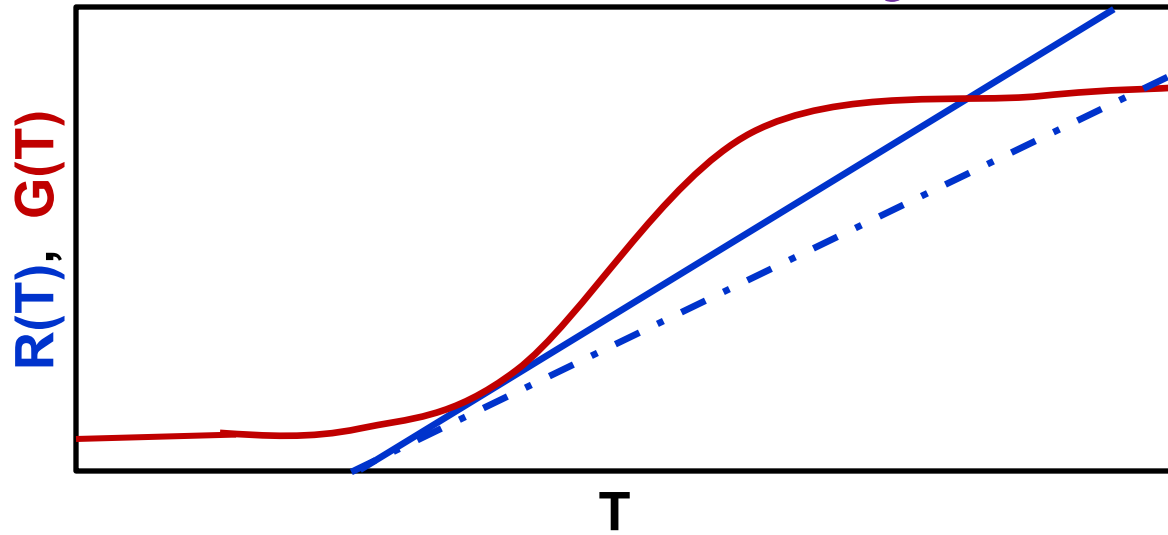
- Slight increase in T above $T_{S,green}$ causes reactor T to jump to $T_{S,cyan}$
- Ignition temp: T where jump from $T_{S,lower}$ to $T_{S,upper}$ occurs
- Slight decrease in T below $T_{S,magenta}$ causes reactor T to drop to $T_{S,yellow}$
- Extinction temp: T where drop from $T_{S,upper}$ to $T_{S,lower}$ occurs



Plot T_S vs T_0



Runaway Reaction



Ignition temperature is very important:
once T_0 exceeds T_{ignition} , transition to the upper steady state will occur

- undesirable
- dangerous

➔ Runaway reaction
 $R(T)$ only intersects with upper steady state

