# CEE 451G ENVIRONMENTAL FLUID MECHANICS LECTURE 1: SCALARS, VECTORS AND TENSORS

A **scalar** has magnitude but no direction.

An example is pressure p.

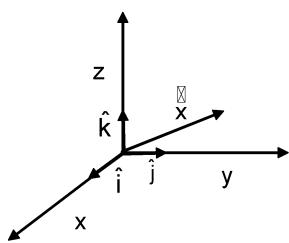
The coordinates x, y and z of Cartesian space are scalars.

A **vector** has both magnitude and direction

Let  $\hat{i}$ ,  $\hat{j}$ ,  $\hat{k}$  denote **unit** vectors in the x, y and z direction. The **hat** denotes a magnitude of unity

The **position vector**  $\overset{\bowtie}{x}$  (the arrow denotes a vector that is not a unit vector) is given as

$$\ddot{x} = x\hat{i} + y\hat{j} + z\hat{k}$$



The **velocity vector**  $\ddot{\mathbf{u}}$  is given as

$$\overset{\boxtimes}{u} = \frac{d\overset{\bowtie}{x}}{dt} = \frac{dx}{dt}\,\hat{i} + \frac{dy}{dt}\,\hat{j} + \frac{dz}{dt}\,\hat{k}$$

The **acceleration** vector  $\ddot{a}$  is given as

$$\mathbf{a} = \frac{d\mathbf{u}}{dt} = \frac{d\mathbf{u}}{dt}\,\mathbf{\hat{i}} + \frac{d\mathbf{v}}{dt}\,\mathbf{\hat{j}} + \frac{d\mathbf{w}}{dt}\,\mathbf{\hat{k}} = \frac{d^2\mathbf{x}}{dt^2} = \frac{d^2\mathbf{x}}{dt^2}\,\mathbf{\hat{i}} + \frac{d^2\mathbf{y}}{dt^2}\,\mathbf{\hat{j}} + \frac{d^2\mathbf{z}}{dt^2}\,\mathbf{\hat{k}}$$

The **units** that we will use in class are length L, time T, mass M and temperature °. The units of a parameter are denoted in brackets. Thus

$$\begin{bmatrix} x \\ x \end{bmatrix} = L$$
 $\begin{bmatrix} u \\ u \end{bmatrix} = LT^{-1}$ 
 $\begin{bmatrix} a \\ a \end{bmatrix} = ?$ 
 $LT^{-2}$ 

**Newton's second law** is a vectorial statement: where F denotes the force vector and m denotes the mass (which is a scalar)

$$F = ma$$

The components of the force vector can be written as follows:

$$\overset{\bowtie}{\mathsf{F}} = \mathsf{F}_{\mathsf{x}}\hat{\mathsf{i}} + \mathsf{F}_{\mathsf{y}}\hat{\mathsf{j}} + \mathsf{F}_{\mathsf{z}}\hat{\mathsf{k}}$$

The **dimensions** of the force vector are the dimension of mass times the dimension acceleration

$$[\overset{\bowtie}{\mathsf{F}}] = [\mathsf{F}_{\mathsf{x}}] = \mathsf{MLT}^{-2}$$

Pressure p, which is a scalar, has dimensions of force per unit area. The dimensions of pressure are thus

$$[p] = MLT^{-2}/(L^2) = ML^{-1}T^{-2}$$

The acceleration of gravity g is a scalar with the dimensions of (of course) acceleration:

$$[g] = LT^{-2}$$

A scalar can be a function of a vector, a vector of a scalar, etc. For example, in fluid flows pressure and velocity are both functions of position and time:

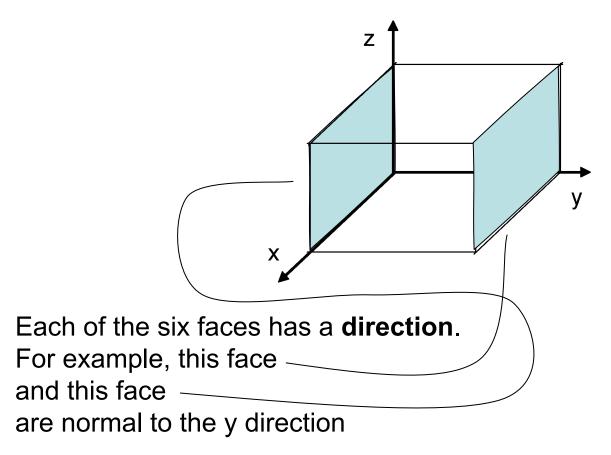
$$p = p(x,t)$$
 ,  $u = u(x,t)$ 

A scalar is a **zero-order tensor**. A vector is a **first-order tensor**. A matrix is a **second order tensor**. For example, consider the **stress tensor T**.

$$oldsymbol{ au} = egin{pmatrix} au_{\mathsf{xx}} & au_{\mathsf{xy}} & au_{\mathsf{xz}} \ au_{\mathsf{yx}} & au_{\mathsf{yy}} & au_{\mathsf{yz}} \ au_{\mathsf{zx}} & au_{\mathsf{zy}} & au_{\mathsf{zz}} \end{pmatrix}$$

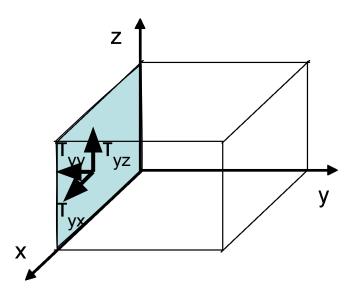
The stress tensor has 9 components. What do they mean? Use the following mnemonic device: **first face**, **second stress** 

Consider the volume element below.



A force acting on any face can act in the x, y and z directions.

Consider the face below.



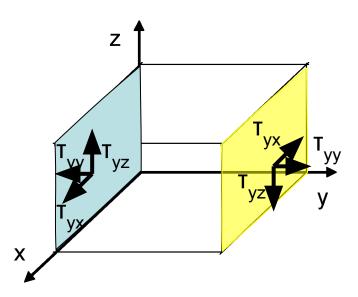
The face is in the direction y.

The force per unit face area acting in the x direction on that face is the stress  $\tau_{yx}$  (first face, second stress).

The forces per unit face area acting in the y and z directions on that face are the stresses  $\tau_{vv}$  and  $\tau_{vz}$ .

Here  $\tau_{yy}$  is a **normal stress** (acts normal, or perpendicular to the face) and  $\tau_{yz}$  are **shear stresses** (act parallel to the face)

Some conventions are in order



Normal stresses are defined to be positive **outward**, so the orientation is reversed on the face located  $\Delta y$  from the origin

Shear stresses similarly reverse sign on the opposite face face are the stresses  $\tau_{vv}$  and  $\tau_{vz}$ .

Thus a positive normal stress puts a body in tension, and a negative normal stress puts the body in compression. Shear stresses always put the body in shear.`

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Another way to write a vector is in **Cartesian** form:

$$\ddot{\mathbf{x}} = \mathbf{x}\hat{\mathbf{i}} + \mathbf{y}\hat{\mathbf{j}} + \mathbf{z}\hat{\mathbf{k}} = (\mathbf{x}, \mathbf{y}, \mathbf{z})$$

The coordinates x, y and z can also be written as  $x_1$ ,  $x_2$ ,  $x_3$ . Thus the vector can be written as

$$\overset{\bowtie}{\mathsf{X}} = (\mathsf{X}_1, \mathsf{X}_2, \mathsf{X}_3)$$

or as

$$\ddot{x} = (x_i), i = 1..3$$

or in **index notation**, simply as

$$\mathbf{X} = \mathbf{X}_{i}$$

where i is understood to be a dummy variable running from 1 to 3.

Thus  $x_i$ ,  $x_j$  and  $x_p$  all refer to the same vector  $(x_1, x_2 \text{ and } x_3)$ , as the index (subscript) always runs from 1 to 3.

**Scalar multiplication**: let  $\alpha$  be a scalar and  $\stackrel{\square}{A} = A_i$  be a vector. Then

$$\alpha \mathbf{A} = \alpha \mathbf{A}_{i} = (\alpha \mathbf{A}_{i}, \alpha \mathbf{A}_{2}, \alpha \mathbf{A}_{3})$$

is a vector.

Dot or scalar product of two vectors results in a scalar:

$$\overset{\bowtie}{\mathsf{A}} \bullet \overset{\bowtie}{\mathsf{B}} = \mathsf{A}_1 \mathsf{B}_1 + \mathsf{A}_2 \mathsf{B}_2 + \mathsf{A}_3 \mathsf{B}_3 = \mathsf{scalar}$$

In index notation, the dot product takes the form

$$\overset{\text{\tiny{$\mathbb{A}$}}}{\mathsf{A}} \bullet \overset{\text{\tiny{$\mathbb{B}$}}}{\mathsf{B}} = \sum_{\mathsf{i}=1}^{3} \mathsf{A}_{\mathsf{i}} \mathsf{B}_{\mathsf{i}} = \sum_{\mathsf{k}=1}^{3} \mathsf{A}_{\mathsf{k}} \mathsf{B}_{\mathsf{k}} = \sum_{\mathsf{r}=1}^{3} \mathsf{A}_{\mathsf{r}} \mathsf{B}_{\mathsf{r}} =$$

Einstein summation convention: if the same index occurs twice, **always** sum over that index. So we abbreviate to

$$A \bullet B = A_i B_i = A_k B_k = A_r B_r$$

There is no free index in the above expressions. Instead the indices are paired (e.g. two i's), implying summation. The result of the dot product is thus a scalar.

Magnitude of a vector:

$$\left| \overset{\boxtimes}{\mathsf{A}} \right|^2 = \overset{\boxtimes}{\mathsf{A}} \bullet \overset{\boxtimes}{\mathsf{A}} = \mathsf{A}_{\mathsf{i}} \mathsf{A}_{\mathsf{i}}$$

A **tensor** can be constructed by multiplying two vectors (not scalar product):

$$A_{i}B_{j} = (A_{i}B_{j}), i = 1..3, j = 1..3 = \begin{pmatrix} A_{1}B_{1} & A_{2}B_{1} & A_{3}B_{1} \\ A_{1}B_{2} & A_{2}B_{2} & A_{3}B_{3} \\ A_{1}B_{3} & A_{2}B_{3} & A_{3}B_{3} \end{pmatrix}$$

Two free indices (i, j) means the result is a **second-order** tensor

Now consider the expression

$$A_iA_iB_i$$

This is a **first-order tensor**, or **vector** because there is only one free index, i (the j's are paired, implying summation).

$$A_i A_j B_j = (A_1 B_1 + A_2 B_2 + A_3 B_2)(A_1, A_2, A_3)$$

That is, scalar times vector = vector.

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# Kronecker delta $\delta_{ij}$

$$\delta_{ij} = \begin{cases} 1 & \text{if} & i = j \\ 0 & \text{if} & i \neq j \end{cases} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Since there are two free indices, the result is a second-order tensor, or matrix. The Kronecker delta corresponds to the **identity matrix**.

Third-order Levi-Civita tensor.

$$\epsilon_{ijl} = \begin{cases} 1 & \text{if} \quad i,j,k \\ -1 & \text{if} \quad i,j,k \\ 0 & \text{otherwise} \end{cases} \text{ cycle clockwise: 1,2,3, 2,3,1 or 3,1,2}$$

Vectorial **cross product**:

$$\overset{\bowtie}{A}x\overset{\bowtie}{B}=\epsilon_{ijk}A_{j}B_{k}$$

One free index, so the result must be a vector.

Vectorial cross product: Let C be given as

$$C = AxB$$

Then

$$(A_2B_3 - A_3B_2)\hat{i} + (A_3B_1 - A_1B_3)\hat{j} + (A_1B_2 - A_2B_1)\hat{k}$$

Vectorial cross product in tensor notation:

$$C_{i} = \epsilon_{ijk} A_{j} B_{k}$$

Thus for example

$$C_1=\epsilon_{1jk}A_jB_k=\epsilon_{123}^\dagger A_2B_3+\epsilon_{132}^\dagger A_3B_2+\epsilon_{111}^\dagger A_1B_1+\qquad \text{a lot of other terms that}\\ =A_2B_3-A_3B_2$$

i.e. the same result as the other slide. The same results are also obtained for  $C_2$  and  $C_3$ .

The **nabla vector operator**  $\stackrel{\bowtie}{\nabla}$  :

$$\nabla = \hat{\mathbf{i}} \frac{\partial}{\partial \mathbf{x}_1} + \hat{\mathbf{j}} \frac{\partial}{\partial \mathbf{x}_2} + \hat{\mathbf{k}} \frac{\partial}{\partial \mathbf{x}_3}$$

or in index notation

$$\frac{\partial}{\partial \mathbf{X_i}}$$

The **gradient** converts a scalar to a vector. For example, where p is pressure,

grad(p) = 
$$\nabla p = \frac{\partial p}{\partial x_1} \hat{i} + \frac{\partial p}{\partial x_2} \hat{j} + \frac{\partial p}{\partial x_3} \hat{k}$$

or in index notation

grad(p) = 
$$\frac{\partial p}{\partial x_i}$$

The single free index i (free in that it is not paired with another i) in the above expression means that grad(p) is a vector.

The **divergence** converts a vector into a scalar. For example, where  $\ddot{u}$  is the velocity vector,

$$\operatorname{div}(\mathbf{u}) = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} = \frac{\partial u_i}{\partial x_i} = \frac{\partial u_k}{\partial x_k}$$

Note that there is no free index (two i's or two k's), so the result is a scalar.

The **curl** converts a vector to a vector. For example, where  $\ddot{\mathbf{u}}$  is the velocity vector,

$$\mathbf{curl}(\mathbf{u}) = \nabla \mathbf{x} \mathbf{u} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial \mathbf{x}_{1}} & \frac{\partial}{\partial \mathbf{x}_{2}} & \frac{\partial}{\partial \mathbf{x}_{3}} \\ \mathbf{u}_{1} & \mathbf{u}_{2} & \mathbf{u}_{3} \end{vmatrix} = \\ \left( \frac{\partial \mathbf{u}_{3}}{\partial \mathbf{x}_{2}} - \frac{\partial \mathbf{u}_{2}}{\partial \mathbf{x}_{3}} \right) \hat{\mathbf{i}} + \left( \frac{\partial \mathbf{u}_{1}}{\partial \mathbf{x}_{3}} - \frac{\partial \mathbf{u}_{3}}{\partial \mathbf{x}_{1}} \right) \hat{\mathbf{j}} + \left( \frac{\partial \mathbf{u}_{2}}{\partial \mathbf{x}_{1}} - \frac{\partial \mathbf{u}_{1}}{\partial \mathbf{x}_{2}} \right) \hat{\mathbf{k}}$$

or in index notation,

$$\operatorname{curl}(\mathbf{u}) = \varepsilon_{ijk} \frac{\partial u_k}{\partial x_i}$$

One free index i (the j's and the k's are paired) means that the result is a vector

A useful manipulation in tensor notation can be used to change an index in an expression:

$$\delta_{ij}u_{j}=u_{i}$$

This manipulation works because the Kronecker delta  $\delta_{ij} = 0$  except when i = j, in which case it equals 1.